

ELLIPTIC EQUATIONS OF HIGHER STOCHASTIC ORDER

S. V. LOTOTSKY, B. L. ROZOVSKII, AND X. WAN

ABSTRACT. This paper discusses analytical and numerical issues related to elliptic equations with random coefficients which are generally nonlinear functions of white noise. Singularity issues are avoided by using the Itô-Skorohod calculus to interpret the interactions between the coefficients and the solution. The solution is constructed by means of the Wiener Chaos (Cameron-Martin) expansions. The existence and uniqueness of the solutions are established under rather weak assumptions, the main of which requires only that the expectation of the highest order (differential) operator is a non-degenerate elliptic operator. The deterministic coefficients of the Wiener Chaos expansion of the solution solve a lower-triangular system of linear elliptic equations (the propagator). This structure of the propagator insures linear complexity of the related numerical algorithms. Using the lower triangular structure and linearity of the propagator, the rate of convergence is derived for a spectral/ hp finite element approximation. The results of related numerical experiments are presented.

1. INTRODUCTION

The objective of this paper is to study, both analytically and numerically, linear elliptic equations with random coefficients. The perturbations introduced by the random coefficients can represent actual physical input or uncertainty about the model. In this paper, we limit our considerations to coefficients that are (possibly nonlinear) functions of Gaussian random fields. A popular example of an equation with random coefficients is Darcy equation for pressure with lognormal permeability field (see [8] and the references therein). It can be specified as the following Dirichlet problem:

$$-\nabla(a(\mathbf{x}, \omega) \nabla u(\mathbf{x})) = f(\mathbf{x}), \mathbf{x} \in \mathcal{O} \subset \mathbb{R}^d, u|_{\partial\mathcal{O}} = g(\mathbf{x}), \quad (1.1)$$

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where

$$a(\mathbf{x}, \omega) = \bar{a}(\mathbf{x}) \exp \left\{ \epsilon(\mathbf{x}) - \frac{1}{2} \mathbb{E} |\epsilon(\mathbf{x})|^2 \right\}, \quad (1.2)$$

ω symbolizes “chance”, $\epsilon(\mathbf{x}) = \sum_{k \geq 1} \sigma_k(\mathbf{x}) \xi_k$ is the noise term, $\sigma_k(\mathbf{x})$ are deterministic functions, and $\boldsymbol{\xi} := \{\xi_k(\omega)\}_{k \geq 1}$ is a set of uncorrelated Gaussian random variables with zero mean and unit variance. The force f and the boundary value g as well as the boundary itself are assumed to be non-random. Clearly, $\epsilon(\mathbf{x})$ is a Gaussian random field with zero mean and covariance $Q(\mathbf{x}, \mathbf{y}) = \sum_k \sigma_k(\mathbf{x}) \sigma_k(\mathbf{y})$ and $\bar{a}(\mathbf{x})$ is the mean of $a(\mathbf{x}, \omega)$: $\bar{a}(\mathbf{x}) = \mathbb{E}[a(\mathbf{x}, \omega)]$.

Another typical example is (1.1) with

$$a(\mathbf{x}, \omega) := \bar{a}(\mathbf{x}) + \epsilon(\mathbf{x}, \omega) \quad (1.3)$$

Models (1.1)–(1.3) have been actively investigated recently in the context of uncertainty quantification problems (see e.g [2]). Note that equation (1.1) with the function $a(\mathbf{x}, \omega)$ from (1.3) is ill-posed because $a(\mathbf{x}, \omega)$ can take negative values.

Linear equation with multiplicative noise are usually referred to as bi-linear. These equations are much more challenging than linear equations with additive random perturbations. In fact, equation (1.1) is often too singular to be “physical”. By this reason, certain regularization is in order for each of these two models.

One effective and popular regularization procedure for these models is to replace equation by its Wick product version

$$-\nabla(a(\mathbf{x}, \omega) \diamond \nabla u(\mathbf{x})) = f(\mathbf{x}), \mathbf{x} \in \mathcal{O} \subset \mathbb{R}^d, u|_{\partial \mathcal{O}} = g(\mathbf{x}); \quad (1.4)$$

see [8] for a discussion of the Wick product as a modeling tool. Admittedly, “Wick product” is anything but product, it is rather a stochastic convolution. In fact, Wick product is a version of Malliavin divergence operator. For example, $a(\mathbf{x}, \omega) \diamond \nabla u(\mathbf{x})$ is the action of the Malliavin divergence operator with respect to $a(\mathbf{x}, \omega)$ on the random field $\nabla u(\mathbf{x})$.

The idea of this regularization procedure was championed by K. Itô in his seminal work [9]. Specifically, Itô has replaced the “product” model

$$\dot{u}(t) = \mathbf{a}(u(t)) + \mathbf{b}(u(t)) \cdot \dot{W}(t)$$

by a “stochastic convolution” model

$$\dot{u}(t) = \mathbf{a}(u(t)) + \mathbf{b}(u(t)) \diamond \dot{W}(t),$$

which turned out to be equivalent to the stochastic Itô differential:

$$\mathbf{b}(u(t)) \diamond \dot{W} dt = \mathbf{b}(u(t)) dW(t).$$

In this paper we extend Itô's approach to elliptic stochastic PDEs. Since absence of the time variable prevents elliptic SPDEs from being “causal”, the underlying Itô integrals are replaced by the Skorohod integrals (see [15]). The subsequent analysis of the equation is based on the *Wiener chaos expansion* (WCE) with respect to the Cameron-Martin basis. The Cameron-Martin basis consists of random variables

$$\xi_{\alpha} = \prod_k \frac{H_{\alpha_k}(\xi)}{\sqrt{\alpha_k!}},$$

where $\alpha = (\alpha_1, \alpha_2, \dots)$ is a multi-index with non-negative integer entries and $H_{\alpha_k}(x)$ is the Hermite polynomial of order α_k (see Section 2 for more detail). The WCE solution of the equation is given by the series $u(\mathbf{x}) = \sum_{\alpha \in \mathcal{J}} u_{\alpha}(\mathbf{x}) \xi_{\alpha}$, where $u_{\alpha} = \mathbb{E}[u \xi_{\alpha}]$. One can view the Cameron-Martin expansion as a Fourier expansion that separates random and deterministic components of the equation. We demonstrate that the deterministic components $u_{\alpha}(\mathbf{x})$ are uniquely defined by a lower triangular system of linear *deterministic* elliptic equations (see (3.5) below). We refer to this system as the (uncertainty) *propagator*. Since the basis elements ξ_{α} are given by explicit formulas, the propagator is the key element of the solution.

From the numerical stand point, it is important that the lower triangular (in fact, *bi-diagonal*) structure of the propagator ensures *linear* computational complexity. While it is possible to define an approximate WCE solution for the standard multiplicative model (1.1), the corresponding propagator turns out a full system.

From the statistical point of view, it is important that a solution of the Wick product model (1.4) is an *unbiased* random perturbation of a solution of the deterministic equation (1.5)

$$-\nabla(\mathbb{E}[a(\mathbf{x}, \omega)] \nabla v(\mathbf{x})) = f(\mathbf{x}), \mathbf{x} \in \mathcal{O} \subset \mathbb{R}^d, u|_{\partial \mathcal{O}} = g(\mathbf{x}) \quad (1.5)$$

in that the mean (expectation) of the solution of equation (1.4) is a solution of equation (1.5).

In this paper we investigate elliptic equations of arbitrary deterministic and stochastic order. As usual, by deterministic order we understand the order of the highest differential operator; in particular, equation (1.1) has deterministic order two. The “stochastic order” of the equation is, by definition, the highest order of the Hermite polynomial in the Wiener chaos

expansion of the random coefficients. For equation (1.1) with the coefficient $a(\mathbf{x}, \omega)$ defined by (1.3) the stochastic order is 1; for the same equation with the log-normal coefficient (1.2) this order is infinite. Note that any $a(\mathbf{x}, \omega)$ which is *nonlinear* function of $\boldsymbol{\xi} := \{\xi_k(\omega)\}_{k \geq 1}$ has stochastic order higher than 1.

The stochastic order is an important and also intricate characteristic of the equations in question. For example, analytically, problem (1.1) with Gaussian $a(\mathbf{x}, \omega)$ (stochastic order 1) is much more difficult than the lognormal $a(\mathbf{x}, \omega)$ (stochastic order ∞). One reason for this is that in the Gaussian case problem (1.1) is ill posed (from the classical point of view). However, as we will see later, this ill posedness is irrelevant in our setting. In particular, it is shown in Section 3 that bilinear stochastic elliptic PDEs are uniquely solvable under a few simple conditions. The key is the following:

Assumption A. *The expectation of the highest order (differential) operator is a nondegenerate elliptic operator.*

For example, in the setting (1.1) this means $\mathbb{E}[a(\mathbf{x}, \omega)] \geq c > 0$.

In Section 4 it is shown that under assumptions that are very close to those that guarantee the existence and uniqueness of the solution one can construct effective finite-element approximation algorithms. More specifically, using the lower triangular structure and the linearity of the uncertainty propagator and incorporating the estimates of the operator norms, we provide an a priori error estimate for the convergence of spectral/ hp finite element method. For equations of stochastic order one, this analysis has been carried out in [22]. The same strategy can be replicated to obtain new a priori error estimates for other numerical methods for solving the propagator.

As mentioned above the WCE solution is given by the series

$$u(\mathbf{x}) = \sum_{\boldsymbol{\alpha} \in \mathcal{J}} u_{\boldsymbol{\alpha}}(\mathbf{x}) H_{\boldsymbol{\alpha}}(\boldsymbol{\xi}) / \sqrt{\boldsymbol{\alpha}!}.$$

However, very simple examples (see Section 3) demonstrate that the variance of WCE solutions is typically infinite or, more precisely, that $\mathbb{E}\|u\|^2 = \sum_{\boldsymbol{\alpha} \in \mathcal{J}} \|u_{\boldsymbol{\alpha}}\|^2 = \infty$. Fortunately, the blow-up of the Wick product model is controllable, in that the WCE solution can be effectively rescaled by simple weights $r_{\boldsymbol{\alpha}}$ such that $\sum_{\boldsymbol{\alpha}} r_{\boldsymbol{\alpha}}^2 u_{\boldsymbol{\alpha}}^2 < \infty$ (see Theorem 3.3). We remark that, in the “standard product” setting, the blow-up of the WCE solution is typically much more severe than in the Wick product setting (see Section 5)

There are alternative ways to address bi-linear elliptic SPDEs. The most developed alternative approach is based on Hida's white noise analysis [7]. The white noise approach exploits the built-in set of stochastic spaces, such as Hida or Kondratiev spaces [11, 12], or even larger exponential spaces [16]. The traditional approach [17, 20, 21, etc.] has to select a stochastic space and then to study the largest possible class of equations admitting a solution in that space. The difference of our approach is that we *select* the stochastic space that is in some sense optimal for the particular equation under consideration.

The paper consists of four sections. Section 2 describes the classes of solution spaces for the equations of interest. The theorem about existence and uniqueness of the solution is in Section 3. Numerical analysis of the equation using finite-element approximation, including the proof of convergence, is in Section 4. Finally, Section 5 presents results of numerical experiments.

2. WEIGHTED CHAOS SPACES

Let \mathcal{J} be the set of multi-indices $\alpha = (\alpha_1, \alpha_2, \dots)$ such that $\alpha_k \in \{0, 1, 2, \dots\}$. For $\alpha, \beta \in \mathcal{J}$, we define

$$|\alpha| = \sum_{k \geq 1} \alpha_k, \quad \alpha! = \prod_{k \geq 1} \alpha_k!, \quad \alpha + \beta = (\alpha_1 + \beta_1, \alpha_2 + \beta_2, \dots).$$

By definition, $\alpha > 0$ if $|\alpha| > 0$ and $\beta \leq \alpha$ if

$$\beta_k \leq \alpha_k \text{ for all } k \geq 1.$$

If $\beta \leq \alpha$, then we define

$$\alpha - \beta = (\alpha_1 - \beta_1, \alpha_2 - \beta_2, \dots).$$

We use the following notations for the special multi-indices:

- (1) $(\mathbf{0})$ is the multi-index with all zero entries: $(\mathbf{0})_k = 0$ for all k ;
- (2) $\epsilon(i)$ is the multi-index of length 1 and with the single non-zero entry at position i :

$$\epsilon(i)_k = \begin{cases} 1, & \text{if } k = i; \\ 0, & \text{if } k \neq i. \end{cases}$$

We also use convention $\epsilon(\mathbf{0}) = (\mathbf{0})$.

Given a sequence $q = (q_k, k \geq 1)$ of positive numbers and $\alpha \in \mathcal{J}$, we define

$$q^\alpha = \prod_{k \geq 1} q_k^{\alpha_k}. \tag{2.1}$$

Next, we introduce the following objects:

- (1) $\mathbb{F} = (\Omega, \mathcal{F}, \mathbb{P})$, a complete probability space that is big enough to support countably many independent Gaussian random variables;
- (2) $\boldsymbol{\xi} = \{\xi_k, k \geq 1\}$, i.i.d. standard normal random variables on \mathbb{F} ;
- (3) V , a real separable Hilbert space;
- (4) $H_n = H_n(t)$, $n = 0, 1, 2, \dots$, $t \in \mathbb{R}$, one-dimensional Hermite polynomial of order n :

$$H_n(t) = (-1)^n e^{t^2/2} \frac{d^n}{dt^n} e^{-t^2/2}. \quad (2.2)$$

In particular, $H_0(t) = 1$, $H_1(t) = t$, $H_2(t) = t^2 - 1$, $H_3(t) = t^3 - 3t, \dots$

- (5) Stochastic Hermite polynomials

$$H_{\boldsymbol{\alpha}}(\boldsymbol{\xi}) = \prod_{k \geq 1} H_{\alpha_k}(\xi_k), \quad \boldsymbol{\alpha} \in \mathcal{J}.$$

For example, if $\boldsymbol{\alpha} = (0, 2, 0, 1, 3, 0, 0, \dots)$, then

$$H_{\boldsymbol{\alpha}}(\boldsymbol{\xi}) = (\xi_2^2 - 1) \xi_4 (\xi_3^3 - 3\xi_5).$$

Recall the following result (Cameron and Martin [3]).

Theorem 2.1. *The collection*

$$\left\{ \frac{H_{\boldsymbol{\alpha}}(\boldsymbol{\xi})}{\sqrt{\boldsymbol{\alpha}!}}, \boldsymbol{\alpha} \in \mathcal{J} \right\}$$

is an orthonormal basis in the space of square-integrable random variables that are measurable with respect to the σ -algebra generated by the collection $\boldsymbol{\xi}$.

In what follows, assume that the σ -algebra \mathcal{F} in \mathbb{F} is generated by $\boldsymbol{\xi}$, and denote by $L_2(\mathbb{F}; V)$ the collection of square-integrable V -valued random elements. By Theorem 2.1, every $v \in L_2(\mathbb{F}; V)$ has a unique representation

$$v = \sum_{\boldsymbol{\alpha} \in \mathcal{J}} v_{\boldsymbol{\alpha}} H_{\boldsymbol{\alpha}}(\boldsymbol{\xi}), \quad (2.3)$$

where

$$v_{\boldsymbol{\alpha}} = \frac{\mathbb{E}[v H_{\boldsymbol{\alpha}}(\boldsymbol{\xi})]}{\boldsymbol{\alpha}!}, \quad (2.4)$$

and

$$\mathbb{E}\|v\|_V^2 = \sum_{\boldsymbol{\alpha} \in \mathcal{J}} \boldsymbol{\alpha}! \|v_{\boldsymbol{\alpha}}\|_V^2. \quad (2.5)$$

Then, given a collection $\mathcal{R} = \{r_\alpha, \alpha \in \mathcal{J}\}$ of positive real numbers, we define the space $\mathcal{RL}_2(\mathbb{F}; V)$ as the closure of $L_2(\mathbb{F}; V)$ in the norm

$$\|v\|_{\mathcal{RL}_2(\mathbb{F}; V)}^2 = \sum_{\alpha \in \mathcal{J}} r_\alpha \alpha! \|v_\alpha\|_V^2. \quad (2.6)$$

The space $\mathcal{RL}_2(\mathbb{F}; V)$ is called a **weighted chaos space**. We use the notation

$$\mathcal{RL}_2(\mathbb{F}) = \mathcal{RL}_2(\mathbb{F}; \mathbb{R}).$$

Proposition 2.2. *Let $\{v_\alpha, \alpha \in \mathcal{J}\}$ be a collection of elements from V . Then there exists a collection \mathcal{R} such that $\sum_{\alpha \in \mathcal{J}} v_\alpha H_\alpha(\xi) \in \mathcal{RL}_2(\mathbb{F}; V)$.*

Proof. We use the following result: if $p > 1$, then

$$\sum_{\alpha \in \mathcal{J}} \frac{1}{(2\mathbb{N})^{p\alpha}} < \infty \quad (2.7)$$

where \mathbb{N} is the sequence of positive integers and

$$(2\mathbb{N})^{p\alpha} = \prod_{k \geq 1} (2k)^{p\alpha_k}$$

(see [8, Proposition 7.1]). Then it is enough to take

$$r_\alpha = \frac{1}{\alpha! (2\mathbb{N})^{2\alpha} (1 + \|v_\alpha\|_V^2)}.$$

□

Definition 2.3. *A generalized random element is a collection $\{v_\alpha, \alpha \in \mathcal{J}\}$ of elements from V .*

3. STOCHASTIC ELLIPTIC EQUATIONS: EXISTENCE AND UNIQUENESS OF SOLUTIONS

Let \mathcal{I} be a subset of \mathcal{J} such that $(\mathbf{0}) \in \mathcal{I}$ and let V' be a real separable Hilbert space. Consider a collection $\{A_\alpha, \alpha \in \mathcal{I}\}$ of bounded linear operators from V to V' , such that the operator $A_{(\mathbf{0})}$ has a bounded inverse. The objective of this section is to study the existence and uniqueness of solution of the following stochastic equation

$$\sum_{\beta \in \mathcal{I}} (A_\beta u) \diamond H_\beta(\xi) = f \quad (3.1)$$

where $f \in V'$ is non-random and \diamond denotes the Wick product [8]. In particular,

$$H_\alpha(\xi) \diamond H_\beta(\xi) = H_{\alpha+\beta}(\xi). \quad (3.2)$$

Definition 3.1. Equation (3.1) is said to have finite stochastic order N if

$$\max_{\alpha \in \mathcal{I}} |\alpha| = N.$$

The equation (3.1) is said to have infinite stochastic order if

$$\sup_{\alpha \in \mathcal{I}} |\alpha| = \infty.$$

A special case of (3.1) is

$$Au + \sum_{k \geq 1} M_k u \diamond \xi_k = f; \quad (3.3)$$

this equations is studied in [14]. In (3.3), $A_\alpha = 0$ for $|\alpha| > 1$, and the equation has stochastic order 1:

$$A_{(0)}u + \sum_{k \geq 1} (A_{\epsilon(k)}u) \diamond H_{\epsilon(k)}(\xi) = f. \quad (3.4)$$

In general, (3.1) can be written as

$$A_{(0)}u = f - \sum_{\beta \in \mathcal{I}, |\beta| > 0} (A_\beta u) \diamond H_\beta(\xi)$$

and interpreted as a bi-linear stochastic perturbation of the deterministic equation $A_{(0)}u = f$; this is the reason for assuming that $(0) \in \mathcal{I}$. The use of the Wick product ensures that the perturbation has zero average.

In view of Proposition 2.2 and property (3.2) of the Hermite polynomials, the following definition is natural.

Definition 3.2. A solution of equation (3.1) is a generalized random element such that

$$\begin{aligned} A_{(0)}u_{(0)} &= f, \quad |\alpha| = 0, \\ A_{(0)}u_\alpha &= - \sum_{\substack{\beta \in \mathcal{I} \\ 0 < \beta \leq \alpha}} A_\beta u_{\alpha-\beta}, \quad |\alpha| > 0. \end{aligned} \quad (3.5)$$

Indeed, writing $u = \sum_{\alpha \in \mathcal{J}} u_\alpha H_\alpha(\xi)$, substituting formally into (3.1), and using (3.2), we find

$$\sum_{\alpha \in \mathcal{J}} \sum_{\beta \in \mathcal{I}} A_\beta u_\alpha H_{\alpha+\beta}(\xi) = f,$$

or

$$\sum_{\alpha \in \mathcal{J}} \left(\sum_{\substack{\beta \in \mathcal{I} \\ \beta \leq \alpha}} A_\beta u_{\alpha-\beta} \right) H_\alpha(\xi) = f.$$

Equating the coefficients of every $H_\alpha(\xi)$, we conclude that (3.5) must hold.

The following theorem establishes the existence and uniqueness of solution of equation (3.1).

Theorem 3.3. *Assume that*

- (1) *The operator $A_{(0)}$ has a bounded inverse $A_{(0)}^{-1}$ from V' to V .*
- (2) *Each of the operators $A_{(0)}^{-1}A_{\beta}$, $\beta \in \mathcal{I}$, $|\beta| > 0$, is bounded on V , with the operator norm C_{β} .*
- (3) *There exists a sequence $b = (b_1, b_2, \dots)$ of positive numbers, such that, for all $\beta \in \mathcal{I}$,*

$$C_{\beta} \leq b^{\beta}. \quad (3.6)$$

- (4) *The function f is deterministic and belongs to V' .*

Then

- (1) *equation (3.1) has a unique solution u ;*
- (2) *the solution is an element of the weighted chaos space $\mathcal{RL}_2(\mathbb{F}; V)$ with*

$$r_{\alpha} = \frac{c^{\alpha}}{\alpha!}$$

for some sequence $c = (c_1, c_2, \dots)$ of positive numbers.

Proof. To begin, let us first understand the structure of the system of equations (3.5). Define

$$B_{\beta} = -A_{(0)}^{-1}A_{\beta}, \quad \beta \in \mathcal{I}.$$

Then (3.5) becomes

$$\begin{aligned} u_{(0)} &= A_{(0)}^{-1}f, \quad |\alpha| = 0, \\ u_{\alpha} &= \sum_{\substack{\beta \in \mathcal{I} \\ 0 < \beta \leq \alpha}} B_{\beta} u_{\alpha - \beta}, \quad |\alpha| > 0. \end{aligned} \quad (3.7)$$

Thus, u_{α} with $|\alpha| = n > 0$ is determined by u_{α} with $|\alpha| < n$, which implies existence and uniqueness of solution of (3.5), and hence of (3.1).

In view of (2.7), to prove that the solution is an element of a weighted chaos space with weights r_{α} of the form $r_{\alpha} = c^{\alpha}/\alpha!$ for some sequence $c = (c_1, c_2, \dots)$ of positive numbers, it is enough to show that there exists a sequence of positive numbers $q = (q_k, k \geq 1)$ such that

$$\|u_{\alpha}\|_V \leq q^{\alpha} \|u_{(0)}\|_V; \quad (3.8)$$

then we can take $c_k = (2kq_k)^{-2}$. Thus, to complete the proof, it remain to establish (3.8).

We start with the particular case (3.3). Even though this equation has been studied in [14], the analysis there relies on a closed-form expression for u_{α} and does not easily extend to higher orders. Denote by B_k the operator $A^{-1}M_k$. Then, for $|\alpha| > 0$, (3.7) becomes

$$u_{\alpha} = \sum_{k \geq 1} \mathbf{1}_{\alpha_k > 0} B_k u_{\alpha - \epsilon(k)}. \quad (3.9)$$

We now show by induction that (3.8) holds with

$$q_k = 4k^2 b_k.$$

Since $u_{\epsilon(k)} = B_k u_{(0)}$, we have (3.8) for $|\alpha| = 1$. Now assume that (3.8) holds for all α with $|\alpha| \leq n-1$. Then (3.9) and the triangle inequality imply, for α with $|\alpha| = n$,

$$\begin{aligned} \|u_{\alpha}\|_V &\leq \|u_{(0)}\|_V \sum_{k \geq 1} \mathbf{1}_{\alpha_k > 0} b_k q^{\alpha - \epsilon(k)} \\ &= q^{\alpha} \|u_{(0)}\|_V \sum_{k \geq 1} \frac{\mathbf{1}_{\alpha_k > 0}}{4k^2} \leq q^{\alpha} \|u_{(0)}\|_V \sum_{k=1}^{\infty} \frac{1}{4k^2} < q^{\alpha} \|u_{(0)}\|_V, \end{aligned}$$

because $\sum_{k \geq 1} k^{-2} = (\pi^2/6) < 4$. This completes the proof of (3.8) for equation (3.3).

Let us now consider the general equation (3.1). We will use assumption (3.6) to show that (3.8) holds for $N > 1$ with

$$q_k = C_0 k^2 b_k,$$

where

$$C_0 = 2 \sum_{\alpha \in \mathcal{J}} \frac{1}{(2\mathbb{N})^{2\alpha}}; \quad (3.10)$$

see (2.7). For the curious reader, we note that numerical computations show that $C_0 \approx 3.2$, and simple analytic considerations put the value of C_0 in the range $[3.0, 3.4]$.

As in the case of equation (3.3), the proof of (3.8) is by induction on $|\alpha|$. Let $k_0 = \min\{|\beta| : \beta \in \mathcal{I}\}$. If $|\alpha| \leq k_0$, then $u_{\alpha} = 0$ and so (3.8) trivially holds. If $\alpha \in \mathcal{I}$ and $|\alpha| = k_0$, then (3.5) implies

$$u_{\alpha} = B_{\alpha} u_{(0)}, \quad \|u_{\alpha}\| \leq b^{\alpha} \|u_{(0)}\|_V \leq q^{\alpha} \|u_{(0)}\|_V.$$

Now assume that (3.8) holds for all u_{α} with $|\alpha| \leq n-1$ and fix $|\alpha| = n > k_0$. By (3.7), the triangle inequality, and the induction assumption,

$$\begin{aligned} \|u_{\alpha}\|_V &\leq \sum_{\substack{\beta \in \mathcal{I} \\ 0 < \beta \leq \alpha}} \|B_{\beta} u_{\alpha - \beta}\|_V \leq \|u_{(0)}\|_V \frac{q^{\alpha}}{C_0} \sum_{\beta \in \mathcal{I}} \frac{1}{(2\mathbb{N})^{2\beta}} \\ &\leq \|u_{(0)}\|_V \frac{q^{\alpha}}{C_0} \sum_{\beta \in \mathcal{J}} \frac{1}{(2\mathbb{N})^{2\beta}} \leq q^{\alpha} \|u_{(0)}\|_V, \end{aligned}$$

where the last inequality follows from (3.10). This completes the proof of Theorem 3.3. \square

Below, we illustrate Theorem 3.3 on several examples. To emphasize the main points, we consider a simplified version of equation (3.1), in which

- (1) $V = V' = \mathbb{R}$, so that every operator A_β is multiplication by a real number; then, with no loss of generality, we take $A_{(0)} = 1$.
- (2) $f = 1$.
- (3) the operators A_β are such that $A_\beta \neq 0$ only if $\beta = (n, 0, 0, \dots)$ for some $n = 0, 1, 2, \dots$.

In other words, (3.1) becomes

$$u = 1 + \sum_{k \geq 1} a_k u \diamond H_k(\xi), \quad (3.11)$$

where ξ is a standard Gaussian random variable. The solution u has the form

$$u = \sum_{n \geq 0} u_n H_n(\xi)$$

for some real numbers u_n , and

$$\mathbb{E}|u|^2 = \sum_{n \geq 1} n! u_n^2. \quad (3.12)$$

As a first example, consider the equation

$$u = 1 + u \diamond \xi. \quad (3.13)$$

It is easy to verify that $u = 1 + \sum_{k \geq 1} H_k(\xi)$ is a solution, that is, $u_k = 1$ for all k . By (3.12), $u \notin L_2(\mathbb{F}; \mathbb{R})$.

The coefficients of the solution of (3.13) are bounded. With a simple modification of (3.13), we can get coefficients that grow exponentially fast. Indeed, consider

$$u = 1 + u \diamond \xi + u \diamond (\xi^2 - 1); \quad (3.14)$$

recall that $\xi^2 - 1 = H_2(\xi)$. Then (3.5) implies

$$u_0 = 1, \quad u_1 = 1, \quad u_n = u_{n-1} + u_{n-2}, \quad n \geq 2.$$

Thus, u_n is n th Fibonacci number, and, as $n \rightarrow \infty$, u_n behaves asymptotically as $(1 + \sqrt{5})^n / 2^n$. In other words, we have $u_n \geq r^n$ for some $r > 1$. Again, (3.12) implies $\mathbb{E}u^2 = \infty$ and so $u \notin L_2(\mathbb{F}; \mathbb{R})$.

Examples (3.13) and (3.14) show that, if equation (3.1) has finite stochastic order, then, in general, the solution does not belong to $L_2(\mathbb{F}; V)$. In fact, we are not aware of any elliptic equation with multiplicative finite-order noise and with a square-integrable solution.

Let us consider now the equation

$$u = 1 - \sum_{n \geq 1} \frac{u \diamond H_n(\xi)}{n!}. \quad (3.15)$$

This equation has infinite stochastic order. Direct computations show that

$$u = 1 + \sum_{n \geq 1} \frac{(-1)^n}{n!} H_n(\xi) = e^{-\xi - (1/2)};$$

one can also derive this by re-writing (3.15) as

$$e^{\xi - (1/2)} \diamond u = 1$$

and noticing that $e^{\xi - (1/2)} \diamond e^{-\xi - (1/2)} = 1$. In particular, $\mathbb{E}u^2 = e < \infty$, so $u \in L_2(\mathbb{F}; \mathbb{R})$.

4. NUMERICAL ANALYSIS OF A STOCHASTIC DIRICHLET PROBLEM

Let $\mathcal{O} \in \mathbb{R}^d$ be a bounded, open, connected set with a Lipschitz continuous boundary $\partial\mathcal{O}$. Denote by $\langle \cdot, \cdot \rangle$ the inner product in $L_2(\mathcal{O})$. Consider the following problem:

$$\begin{cases} Au + \sum_{\substack{\beta \in \mathcal{I} \\ |\beta| \geq 1}} A_\beta u \diamond H_\beta(\xi) = f, & \mathbf{x} \in \mathcal{O}, \\ u = 0, & \mathbf{x} \in \partial\mathcal{O}, \end{cases} \quad (4.1)$$

where ξ is a finite- or infinite-dimensional standard Gaussian random variable. For simplicity, we assume that the force term $f(\mathbf{x})$ is a deterministic function. The operators $A = A_{(0)}$ and A_β are defined by

$$Au(\mathbf{x}) = - \sum_{i,j=1}^d D_i(a_{ij}(\mathbf{x}) D_j u(\mathbf{x})), \quad (4.2)$$

$$A_\beta(\mathbf{x}) = w_\beta(\mathbf{x}) \sum_{i,j=1}^d D_i(\sigma_{ij,\beta}(\mathbf{x}) D_j u(\mathbf{x})), \quad (4.3)$$

where the functions $a_{ij}(\mathbf{x})$, $\sigma_{ij,\beta}(\mathbf{x})$, and $w_\beta(\mathbf{x})$ satisfy the following assumptions:

Assumption 4.1.

- (a) *The functions $a_{ij}(\mathbf{x})$ are measurable and bounded in the closure $\bar{\mathcal{O}}$ of \mathcal{O} ;*

- (b) *There exist positive numbers A_1, A_2 such that $A_1|\mathbf{y}|^2 \leq a_{ij}(\mathbf{x})y_i y_j \leq A_2|\mathbf{y}|^2$ for all $\mathbf{x} \in \bar{\mathcal{O}}$ and $\mathbf{y} \in \mathbb{R}^d$;*
- (c) *The functions $\sigma_{ij,\beta}(\mathbf{x})$ are bounded and measurable:*

$$|\sigma_{ij,\beta}(\mathbf{x})| \leq C_{\beta}^{\sigma} \text{ for all } \mathbf{x} \in \bar{\mathcal{O}}, i, j = 1, \dots, d; \quad (4.4)$$

- (d) *The functions $w_{\beta}(\mathbf{x})$ are bounded and Lipschitz continuous:*

$$|w_{\beta}(\mathbf{x})| \leq C_{\beta}^w \text{ and } |w_{\beta}(\mathbf{x}) - w_{\beta}(\mathbf{y})| \leq C_{\beta}^L |\mathbf{x} - \mathbf{y}|, \quad \mathbf{x}, \mathbf{y} \in \bar{\mathcal{O}}. \quad (4.5)$$

- (e) *There exists a sequence $b = (b_1, b_2, \dots)$ of positive numbers such that, for all $\beta \in \mathcal{I}$,*

$$(C_{\beta}^L + C_{\beta}^w)C_{\beta}^{\sigma} \leq b^{\beta}.$$

Taking as V and V' the Sobolev spaces $H_0^1(\mathcal{O})$ and $H^{-1}(\mathcal{O})$, respectively, we note that Theorem 3.3 applies to equation (4.1) because the operators $A^{-1}A_{\beta}$ are bounded on V . Indeed, note that

$$\|u\|_V^2 = \|u\|_{L_2(\mathcal{O})}^2 + \|\nabla u\|_{L_2(\mathcal{O})}^2$$

and let C_{π} be the Poincare constant for the domain \mathcal{O} :

$$\|u\|_{L_2(\mathcal{O})}^2 \leq C_{\pi} \|\nabla u\|_{L_2(\mathcal{O})}^2. \quad (4.6)$$

Assumption 4.1(b) implies

$$\langle Au, u \rangle \geq \frac{A_1}{1 + C_{\pi}} \|u\|_V^2, \quad u \in V;$$

here $\langle \cdot, \cdot \rangle$ denotes the duality pairing between V' and V relative to the inner product in $L_2(\mathcal{O})$. Then, by direct computation,

$$\|A^{-1}A_{\beta}u\|_V \leq \left(\frac{1 + C_{\pi}}{A_1} \right) (C_{\beta}^L + C_{\beta}^w)C_{\beta}^{\sigma}.$$

Example 4.2. *Let $G(\mathbf{x})$ is a Gaussian field on \mathcal{O} with representation*

$$G(\mathbf{x}) = \sum_{k=1}^{\infty} h_k(\mathbf{x})\xi_k, \quad (4.7)$$

where $\{h_k, k \geq 1\}$ are smooth bounded functions such that

$$\sum_{k \geq 1} \left(\sup_{\mathbf{x} \in \mathcal{O}} |h_k(\mathbf{x})|^2 + \sup_{\mathbf{x} \in \mathcal{O}} |\nabla h_k(\mathbf{x})|^2 \right) < \infty \quad (4.8)$$

and $\xi = \{\xi_k, k \geq 1\}$ is a collection of independent standard Gaussian random variables.

Note that

$$\bar{G}(\mathbf{x}) := \mathbb{E} e^{G(\mathbf{x})} = \exp \left(\frac{1}{2} \sum_{k \geq 1} h_k^2(\mathbf{x}) \right).$$

Examples of (4.7) are Karhuan-Loève expansion [13], spectral expansion [19], and Fourier-Wavelet expansion [6].

Consider the equation

$$\nabla \cdot (e^{G(\mathbf{x})} \diamond \nabla u) = f. \quad (4.9)$$

If

$$\mathfrak{h}_{\alpha}(\mathbf{x}) = \frac{\bar{G}(\mathbf{x})}{\alpha!} \prod_k h_k^{\alpha_k}(\mathbf{x}), \quad \alpha \in \mathcal{J},$$

then

$$e^{G(\mathbf{x})} = \bar{G}(\mathbf{x}) \left(1 + \sum_{|\beta| \geq 1} \mathfrak{h}_{\beta}(\mathbf{x}) \mathbf{H}_{\beta}(\xi) \right).$$

and (4.9) becomes a particular case of (4.1):

$$\nabla \cdot \left(\bar{G}(\mathbf{x}) \nabla u \right) + \sum_{|\beta| \geq 1} \nabla \cdot \left(\bar{G}(\mathbf{x}) \mathfrak{h}_{\beta}(\mathbf{x}) \nabla u \right) \diamond \mathbf{H}_{\beta}(\xi) = f;$$

to satisfy Assumption 4.1(e), take $b_k = \sup_{\mathbf{x} \in \mathcal{O}} |h_k(\mathbf{x})| + \sup_{\mathbf{x} \in \mathcal{O}} |\nabla h_k(\mathbf{x})|$.

Replacing in (4.9) the usual exponential $e^{G(\mathbf{x})}$ with the Wick exponential

$$e^{\diamond G(\mathbf{x})} = e^{G(\mathbf{x}) - \frac{1}{2} \sum_{k \geq 1} h_k^2(\mathbf{x})} \quad (4.10)$$

also leads to an equation of the form (4.1).

The objective of this section is to study numerical solution of (4.1). Specifically, we construct an approximation of (4.1) using finite elements in the physical space \mathcal{O} and a truncated chaos expansion in the probability space. For the sake of concreteness we assume that the physical space is two-dimensional: $d = 2$.

Let \mathcal{T}_h be a family of triangulations of \mathcal{O} with straight edges and h the maximum size of the elements in \mathcal{T}_h . We assume that the family is regular, in other words, the minimal angle of all the triangles is bounded from below by a positive constant. We define the finite element space as

$$V_{h,p}^K = \left\{ v \mid v \circ F_K^{-1} \in \mathcal{P}_p(R) \right\}, \quad V_{h,p} = \left\{ v \in H^1(D) \mid v|_K \in V_{h,p}^K, K \in \mathcal{T}_h \right\},$$

where F_K is the mapping function for the element K which maps the reference element R (for example, an equilateral triangle or an isosceles right triangle) to the element K and $\mathcal{P}_p(R)$ denotes the set of polynomials of degree at most p on R . We assume that $v_h|_{\partial\mathcal{O}} = 0$, $v_h \in V_{h,p}$. Thus, $V_{h,p}$ is an approximation of $H_0^1(\mathcal{O})$ by piece-wise polynomial

functions. There exist many choices of basis functions on the reference elements, such as h -type finite elements [5], spectral/ hp elements [10, 18], etc.

To describe truncated chaos expansion, fix positive integers M, n and define the finite dimensional set $\mathcal{J}_{M,n} \subset \mathcal{J}$ as

$$\mathcal{J}_{M,n} := \left\{ \boldsymbol{\alpha} \mid |\boldsymbol{\alpha}| \leq M, \alpha_k = 0, k > n \right\}. \quad (4.11)$$

The truncated Wiener chaos space $V^{M,n}$ is then

$$V^{M,n} := \left\{ \sum_{\boldsymbol{\alpha} \in \mathcal{J}_{M,n}} f_{\boldsymbol{\alpha}} H_{\boldsymbol{\alpha}}(\boldsymbol{\xi}) \mid f_{\boldsymbol{\alpha}} \in \mathbb{R} \right\}, \quad (4.12)$$

Similarly, let

$$\mathcal{I}_{M,n} = \mathcal{I} \cap \mathcal{J}_{M,n}.$$

It follows from (3.5) that $u_{\boldsymbol{\alpha}}$ is fully determined by coefficients $u_{\boldsymbol{\beta}}$ and operators $A_{\boldsymbol{\beta}}$ with $\boldsymbol{\beta} \leq \boldsymbol{\alpha}$. Thus we only need to keep the operators $A_{\boldsymbol{\beta}}$ with $\boldsymbol{\beta} \in \mathcal{I}_{M,n}$ if we only consider the coefficients $u_{\boldsymbol{\alpha}}$ with $\boldsymbol{\alpha} \in \mathcal{J}_{M,n}$.

The sFEM (stochastic finite element method) can be formulated as follows: Find $u_{h,p}^{M,n} \in V_{h,p} \otimes V^{M,n}$ such that, for all $v \in V_{h,p} \otimes V^{M,n}$,

$$\mathbb{E} \langle A_{(0)} u_{h,p}^{M,n}, v \rangle + \mathbb{E} \sum_{\substack{\boldsymbol{\beta} \in \mathcal{I}_{M,n} \\ |\boldsymbol{\beta}| > 0}} \langle A_{\boldsymbol{\beta}} u_{h,p}^{M,n} \diamond H_{\boldsymbol{\beta}}(\boldsymbol{\xi}), v \rangle = \mathbb{E} \langle f, v \rangle. \quad (4.13)$$

Remark 4.3. *The underlying idea of sFEM is to consider a finite element approximation, followed by a truncation of the chaos expansion. Alternatively, it is possible to truncate the chaos expansion first and then apply the finite element method to the truncated version of the uncertainty propagator (3.5):*

$$\begin{aligned} A_{(0)} u_{(0)} &= f, \quad |\boldsymbol{\alpha}| = 0, \\ A_{(0)} u_{\boldsymbol{\alpha}} + \sum_{\substack{\boldsymbol{\beta} \in \mathcal{I}_{M,n} \\ 0 < \boldsymbol{\beta} \leq \boldsymbol{\alpha}}} A_{\boldsymbol{\beta}} u_{\boldsymbol{\alpha} - \boldsymbol{\beta}} &= 0, \quad |\boldsymbol{\alpha}| > 0. \end{aligned} \quad (4.14)$$

This approach will not be discussed in this paper.

The following is the main result about the convergence of the stochastic finite element method.

Theorem 4.4. *Assume that conditions given in Assumption 4.1 and Theorem 3.3 are satisfied and in addition, $u \in \mathcal{RL}_2(\mathbb{F}; H^{m+1}(\mathcal{O}))$ for some integer $m \geq 1$. Then*

$$\begin{aligned} \|u - u_{h,p}^{M,n}\|_{\mathcal{RL}_2(\mathbb{F}; H_0^1(\mathcal{O}))} &\leq C_1 h^{\mu-1} p^{-m} \|u\|_{\mathcal{RL}_2(\mathbb{F}; H^{m+1}(\mathcal{O}))} \\ &\quad + C_2 \|f\|_{H^{-1}(\mathcal{O})} \left(\frac{1}{2^M} + \frac{1}{\sqrt{n}} \right), \end{aligned} \quad (4.15)$$

where $\mu = \min(m+1, p+1)$, the constants C_1 and C_2 are independent of h and p , and C_2 is also independent of M, n .

Proof. Define $V = H_0^1(\mathcal{O})$ and let the Wiener-chaos expansion of u and $u_{h,p}^{M,n}$ be

$$u = \sum_{\alpha \in \mathcal{J}} u_\alpha H_\alpha(\xi), \quad u_{h,p}^{M,n} = \sum_{\alpha \in \mathcal{J}_{M,n}} \hat{u}_\alpha H_\alpha(\xi),$$

respectively. Then

$$u - u_{h,p}^{M,n} = \sum_{\alpha \in \mathcal{J}_{M,n}} (u_\alpha - \hat{u}_\alpha) H_\alpha + \sum_{\alpha \in \mathcal{J} \setminus \mathcal{J}_{M,n}} u_\alpha H_\alpha, \quad (4.16)$$

and

$$\begin{aligned} \|u - u_{h,p}^{M,n}\|_{\mathcal{RL}_2(\mathbb{F}; V)}^2 &= \sum_{\alpha \in \mathcal{J}_{M,n}} \|u_\alpha - \hat{u}_\alpha\|_V^2 \|H_\alpha\|_{\mathcal{RL}_2(\mathbb{F})}^2 \\ &\quad + \sum_{\alpha \in \mathcal{J} \setminus \mathcal{J}_{M,n}} \|u_\alpha\|_V^2 \|H_\alpha\|_{\mathcal{RL}_2(\mathbb{F})}^2 = I_1 + I_2. \end{aligned} \quad (4.17)$$

The structure of equation (3.5) ensures that the coefficients u_α with $\alpha \in \mathcal{J} \setminus \mathcal{J}_{M,n}$ and operators A_β with $\beta \in \mathcal{I} \setminus \mathcal{I}_{M,n}$ do not affect \hat{u}_α with $\alpha \in \mathcal{J}_{M,n}$. In other words, I_1 is the error from the finite element discretization and I_2 is from the truncation of the chaos expansion.

We first look at the term I_1 by considering equation (4.14). Since the equation for $u_{(0)}$ is a usual deterministic elliptic equation, we here only consider the equations for u_α with $|\alpha| > 0$. By (4.13),

$$A_{(0)} \hat{u}_\alpha + \sum_{\substack{\beta \in \mathcal{I}_{M,n} \\ 0 < \beta \leq \alpha}} A_\beta \hat{u}_{\alpha-\beta} = 0. \quad (4.18)$$

Furthermore, although we truncate the number of random variables and the multi-index set \mathcal{I} , the structure of the uncertainty propagator (3.5) ensures that the exact solution also satisfied

$$A_{(0)} u_\alpha + \sum_{\substack{\beta \in \mathcal{I}_{M,n} \\ 0 < \beta \leq \alpha}} A_\beta u_{\alpha-\beta} = 0. \quad (4.19)$$

Denote by $a(\cdot, \cdot)$ the bilinear form generated by the elliptic operator $A_{(\mathbf{0})}$:

$$a(u, v) = \langle Au, v \rangle.$$

Then, for every $v_h \in V_{h,p}$,

$$\begin{aligned} & a(\hat{u}_\alpha - v_h, \hat{u}_\alpha - v_h) \\ = & a(u_\alpha - v_h, \hat{u}_\alpha - v_h) + a(\hat{u}_\alpha - u_\alpha, \hat{u}_\alpha - v_h) \\ = & a(u_\alpha - v_h, \hat{u}_\alpha - v_h) - \sum_{\substack{\beta \in \mathcal{I}_{M,n} \\ 0 < \beta \leq \alpha}} \langle A_\beta(u_{\alpha-\beta} - \hat{u}_{\alpha-\beta}), \hat{u}_\alpha - v_h \rangle. \end{aligned} \quad (4.20)$$

Assumption 4.1(b) implies

$$\frac{A_1}{1 + C_\pi} \|u\|_V^2 \leq a(u, u), \quad a(u, v) \leq A_2 \|u\|_V \|v\|_V, \quad u \in V,$$

where C_π is the Poincare constant for the domain \mathcal{O} ; see (4.6). Then, with $C_1 = A_1/(1+C_\pi)$ and thinking $w_h = \hat{u}_\alpha - v_h$,

$$C_1 \|\hat{u}_\alpha - v_h\|_V \leq A_2 \|u_\alpha - v_h\|_V + \sup_{\substack{w_h \in V_{h,p} \\ w_h \neq 0}} \frac{\left| \sum_{\substack{\beta \in \mathcal{I}_{M,n} \\ 0 < \beta \leq \alpha}} \langle A_\beta(u_{\alpha-\beta} - \hat{u}_{\alpha-\beta}), w_h \rangle \right|}{\|w_h\|_V}.$$

Combining the above inequality with the triangle inequality

$$\|u_\alpha - \hat{u}_\alpha\|_V \leq \|u_\alpha - v_h\|_V + \|\hat{u}_\alpha - v_h\|_V,$$

we find

$$\begin{aligned} \|u_\alpha - \hat{u}_\alpha\|_V \leq & \left(1 + \frac{A_2}{C_1}\right) \inf_{v_h \in V_{h,p}} \|u_\alpha - v_h\|_V \\ & + \frac{1}{C_1} \sup_{\substack{w_h \in V_{h,p} \\ w_h \neq 0}} \frac{\left| \sum_{\substack{\beta \in \mathcal{I}_{M,n} \\ 0 < \beta \leq \alpha}} \langle A_\beta(u_{\alpha-\beta} - \hat{u}_{\alpha-\beta}), w_h \rangle \right|}{\|w_h\|_V}. \end{aligned} \quad (4.21)$$

By Assumptions 4.1(c,d),

$$\langle A_\beta u, v \rangle \leq C_\beta \|u\|_V \|v\|_V \quad (4.22)$$

where $C_\beta = C_\beta^\sigma(C_\beta^L + C_\beta^w)$. Therefore,

$$\|u_\alpha - \hat{u}_\alpha\|_V \leq C \inf_{v_h \in V_{h,p}} \|u_\alpha - v_h\|_V + \sum_{\substack{\beta \in \mathcal{I}_{M,n} \\ 0 < \beta \leq \alpha}} \hat{C}_\beta \|u_{\alpha-\beta} - \hat{u}_{\alpha-\beta}\|_V \quad (4.23)$$

for some constants C and \hat{C}_β . The first term on the right-hand side is the approximation error of finite element approximation and the second term is the error propagated from the approximation of lower-order coefficients. Using equation (4.23) recursively, we obtain

$$\|u_\alpha - \hat{u}_\alpha\|_V \leq \sum_{\beta \leq \alpha} \hat{C}_{\alpha,\beta} \inf_{v_h \in V_{h,p}} \|u_\beta - v_h\|_V, \quad (4.24)$$

with some constants $\hat{C}_{\alpha,\beta}$. According to the theory of the finite element approximation [1, 5], there exists a positive number C_f such that, for every function $u \in H_0^1(\mathcal{O}) \cap H^{m+1}(\mathcal{O})$,

$$\inf_{v_h \in V_h} \|u_\alpha - v_h\|_V \leq C_f h^{\mu-1} p^{-m} \|u_\alpha\|_{H^{m+1}(D)}, \quad (4.25)$$

where $\mu = \min(m+1, p+1)$. Then

$$\begin{aligned} \|u_\alpha - \hat{u}_\alpha\|_V^2 &\leq C_f^2 h^{2\mu-2} p^{-2m} \left(\sum_{\beta \leq \alpha} \hat{C}_{\alpha,\beta} \|u_\alpha\|_{H^{m+1}(\mathcal{O})} \right)^2 \\ &\leq C_f^2 h^{2\mu-2} p^{-2m} \sum_{\beta \leq \alpha} \hat{C}_{\alpha,\beta}^2 \sum_{\beta \leq \alpha} \|u_\beta\|_{H^{m+1}(\mathcal{O})}^2 \\ &= h^{2\mu-2} p^{-2m} \hat{C}_\alpha \sum_{\beta \leq \alpha} \|u_\beta\|_{H^{m+1}(\mathcal{O})}^2, \end{aligned}$$

and the term I_1 can be bounded as

$$\begin{aligned} I_1 &\leq h^{2\mu-2} p^{-2m} \sum_{\alpha \in \mathcal{J}_{M,n}} \hat{C}_\alpha \sum_{\beta \leq \alpha} \|u_\beta\|_{H^{m+1}(D)}^2 \|H_\alpha\|_{\mathcal{RL}_2(\mathbb{F})}^2 \\ &\leq h^{2\mu-2} p^{-2m} C(M, n) \sum_{\alpha \in \mathcal{J}_{M,n}} \hat{C}_\alpha \|u_\alpha\|_{H^{m+1}(D)}^2 \|H_\alpha\|_{\mathcal{RL}_2(\mathbb{F})}^2 \\ &\leq h^{2\mu-2} p^{-2m} C(M, n) \max_{\alpha \in \mathcal{J}_{M,n}} \hat{C}_\alpha \sum_{\alpha \in \mathcal{J}_{M,n}} \|u_\alpha\|_{H^{m+1}(D)}^2 \|H_\alpha\|_{\mathcal{RL}_2(\mathbb{F})}^2 \\ &\leq h^{2\mu-2} p^{-2m} C_1(M, n) \|u\|_{\mathcal{RL}_2(\mathbb{F}; H^{m+1}(D))}^2. \end{aligned} \quad (4.26)$$

To estimate the truncation error I_2 of the Wiener chaos expansion, note that $\|H_\alpha\|_{\mathcal{RL}_2(\mathbb{F})}^2 = r_\alpha \alpha!$, and, by (3.8),

$$r_\alpha = \frac{(2\mathbb{N})^{-2\alpha} q^{-\alpha}}{\alpha!}.$$

Also,

$$\|u_\alpha\|_V \leq q^\alpha \|u_{(0)}\|_V \leq C_A q^\alpha \|f\|_{H^{-1}(D)}.$$

As a result,

$$I_2 = \sum_{\alpha \in \mathcal{J} \setminus \mathcal{J}_{M,n}} \|u_\alpha\|_V^2 \|H_\alpha\|_{\mathcal{R}L_2(\mathbb{F})}^2 \leq C_A \|f\|_{H^{-1}(D)} \sum_{\alpha \in \mathcal{J} \setminus \mathcal{J}_{M,n}} (2\mathbb{N})^{-2\alpha} \leq C_2 \left(\frac{1}{2^{2M}} + \frac{1}{n} \right),$$

where the last inequality is taken from [4, proof of Theorem 2]. This, together with (4.17) and (4.26), completes the proof of Theorem 4.4. \square

5. COMPUTATIONS

In this section we present some results of numerical experiments and compare our model with another one studied in the literature on uncertainty quantification. We consider the following one-dimensional problem

$$\begin{cases} -\frac{d}{dx}(K * \frac{d}{dx}u) = 1, \\ u(0) = 0, \quad u(1) = 0, \end{cases} \quad (5.1)$$

where

- $*$ denotes either the Wick product \diamond or the ordinary product \cdot ;
- $K = e^{c\xi - \frac{1}{2}c^2}$, ξ is a standard Gaussian random variable, $c > 0$ is a real number.

The random variable K can be regarded as a simplified version of the random field $e^{\diamond G(x)}$ from (4.10). Note that

$$\mathbb{E}[K] = 1, \quad \text{Var}[K] = e^{c^2} - 1. \quad (5.2)$$

Denote by \bar{H}_i the normalized Hermite polynomials:

$$\bar{H}_i(\xi) = \frac{H_i(\xi)}{\sqrt{i!}}.$$

Substituting the truncated Wiener chaos expansion

$$u(x) = \sum_{i=0}^n u_i(x) \bar{H}_i(\xi)$$

into equation (5.1) and implementing the finite element projection we obtain the uncertainty propagators

$$-\sum_{i=0}^n \frac{d^2}{dx^2} u_i(x) \mathbb{E}[K \bar{H}_i \bar{H}_j] = \delta_{0,j}, \quad j = 0, 1, \dots, n \quad (5.3a)$$

$$-\sum_{i=0}^n \frac{d^2}{dx^2} u_i(x) \mathbb{E}[(K \diamond \bar{H}_i) \bar{H}_j] = \delta_{0,j}, \quad j = 0, 1, \dots, n \quad (5.3b)$$

for the ordinary product and the Wick product, respectively.

Using the Taylor expansion and the properties of Hermite polynomials, we obtain

$$\mathbb{E} [K \bar{H}_i \bar{H}_j] = e^{-c^2/2} \sum_{n=0}^{\infty} \sum_{k=0}^{[n/2]} \sum_{p \leq i \wedge j} A(n, k) B(i, j, p) \delta_{n-2k, i+j-2p}, \quad (5.4a)$$

$$\mathbb{E} [(K \diamond \bar{H}_i) \bar{H}_j] = \frac{e^{-c^2/2}}{\sqrt{i!}} \sum_{n=0}^{\infty} \sum_{k=0}^{[n/2]} c^n \frac{\sqrt{(n-2k+i)!}}{(n-2k)!(2k)!} \delta_{n-2k+i, j}, \quad (5.4b)$$

where

$$A(n, k) = \frac{c^n}{n!} \binom{n}{2k} (2k-1)!! \sqrt{(n-2k)!}, \quad B(i, j, p) = \left[\binom{i}{p} \binom{j}{p} \binom{i+j-2p}{i-p} \right]^{1/2}.$$

Formulas similar to (5.4a) and (5.4b) can also be derived when

$$K(x) = e^{\diamond G(x)}, \quad (5.5)$$

with $G(x)$ given by (4.7). For simplicity, we present numerical results for the space-independent case $G(x) = c\xi$, corresponding to $K = e^{c\xi - (c^2/2)}$. The qualitative behavior of the solutions of (5.1), with either ordinary or Wick product, remains the same for the more general noise (5.5), as long as conditions (4.8) hold.

Let us compare the matrix structures of $\mathbb{E} [K \bar{H}_i \bar{H}_j]$ and $\mathbb{E} [(K \diamond \bar{H}_i) \bar{H}_j]$.

• **Ordinary product.** For all $c_1, \dots, c_n \in \mathbb{R}$, we have

$$\begin{aligned} \sum_{i,j=1}^n c_i c_j \mathbb{E} [K \bar{H}_i \bar{H}_j] &= \mathbb{E} \left[\sum_{i,j=1}^n c_i c_j K \bar{H}_i \bar{H}_j \right] \\ &= \mathbb{E} \left[\left(\sum_{i=1}^n K^{1/2} \bar{H}_i c_i \right)^2 \right] \geq 0. \end{aligned}$$

It can be shown that when K is lognormal, c_i must be zero if the equality is satisfied. In other words, the matrix $\mathbb{E} [K \bar{H}_i \bar{H}_j]$ is symmetric and positive definite. Thus all the eigenvalues of the matrix $(\mathbb{E} [K \bar{H}_i \bar{H}_j], i, j = 1, \dots, n)$ are positive, and the system (5.3a) is well-posed. However, all unknowns in equation (5.3a) are coupled together, which is not a desirable property for numerical computation.

- **Wick product.** As n must be even for $i = j$ to get nonzero terms in equation (5.4b), it follows that $\mathbb{E}[(K \diamond \bar{H}_i)\bar{H}_i] = 1$ for all $i = 0, 1, \dots$. Indeed,

$$\begin{aligned} \mathbb{E}[(K \diamond \bar{H}_i)\bar{H}_i] &= \frac{e^{-1/2}}{\sqrt{i!}} \sum_{m=0}^{\infty} \sum_{k=0}^m \frac{\sqrt{(2m-2k+i)!}}{(2m-2k)!(2k)!!} \delta_{2m-2k+i,i} \\ &= \frac{e^{-1/2}}{\sqrt{i!}} \sum_{m=0}^{\infty} \frac{\sqrt{i!}}{(2m)!!} \\ &= e^{-1/2} \sum_{m=0}^{\infty} \frac{1}{2^m m!} = e^{-1/2} e^{1/2} = 1. \end{aligned}$$

Furthermore, we note that $\mathbb{E}[(K \diamond \bar{H}_i)\bar{H}_j]$ is a lower-triangular matrix. In other words, the coefficient u_i depends only on u_j with $j < i$.

In the first numerical experiment, we take $c = 1$ and compare the chaos coefficients of solutions corresponding to the ordinary and Wick products. Figure 1 presents the results. For a better comparison, we use the same scale on both graphs and also plot the function U_d , the solution of

$$-\frac{d^2 U_d(x)}{dx^2} = 1, \quad 0 < x < 1; \quad U_d(0) = U_d(1) = 0, \quad (5.6)$$

Note that U_d is the solution of (5.1) when K is replaced by its mean value $\mathbb{E}[K] = 1$.

We make the following observations:

- (1) When the ordinary product is used, $U_d \neq u_0$; when the Wick product is used, $U_d = u_0$. This is a direct consequence of equations (5.3a) and (5.3b). In Figure 2, we plot the function u_0 corresponding to the ordinary product, for different values of c . It appears that u_0 is an increasing function of c and is close to U_d when $c \leq 0.1$.
- (2) The larger the number i , the smaller the coefficient u_i . However, it appears that the coefficients for the Wick product are much smaller than the corresponding coefficients for the ordinary product. Below, we investigate this phenomenon further by changing the value of c .

In the second numerical experiment, we study the relation between the variance of the perturbation K and the variance of the corresponding solution. We change the value of c from 0.1 to 2.5, which, by (5.2), is equivalent to changing the standard deviation of K from 0.1003 to 22.7379. Figures 3 and 4 present the results.

In Figure 3, we plot the variance of the solution as a function of x for several values of c , and notice that, in each case, the variance seems to be largest near the middle of the

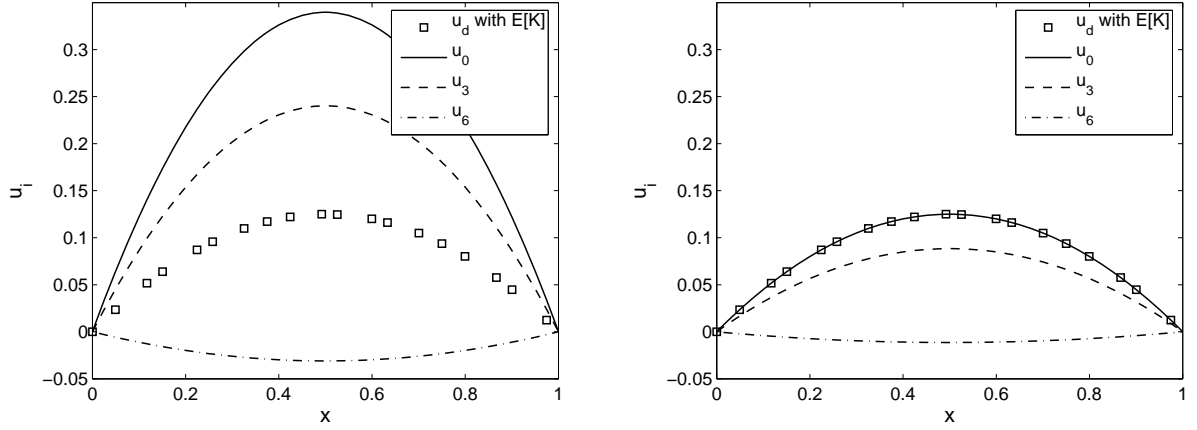
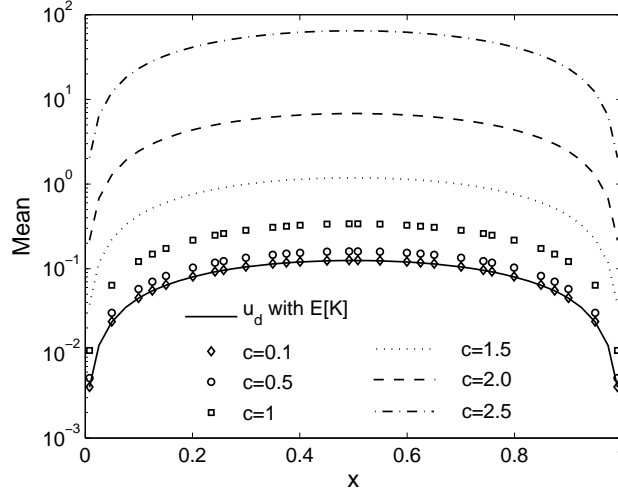
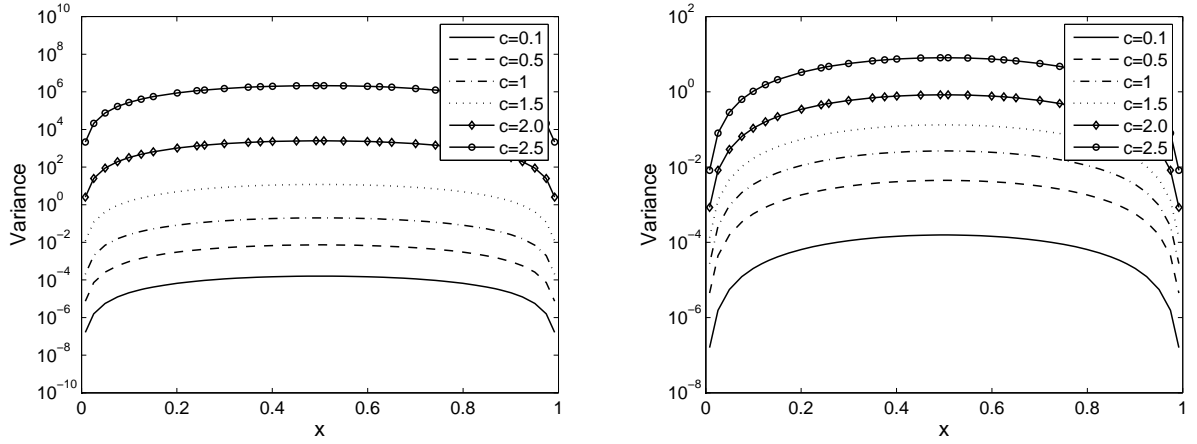


FIGURE 1. Chaos expansion coefficients. Left: ordinary product; Right: Wick product.

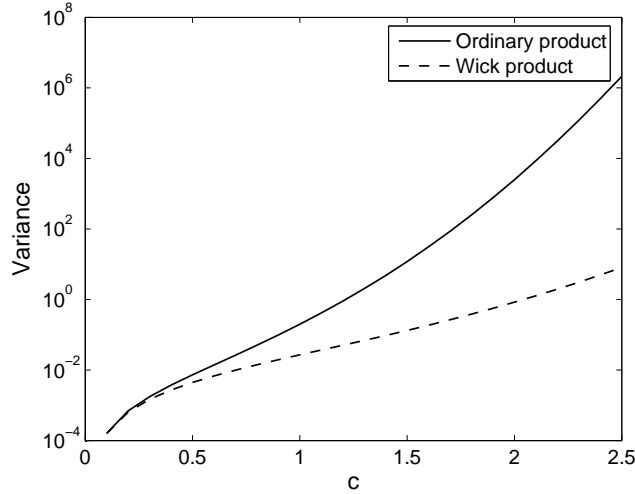
FIGURE 2. The plot of $\mathbb{E}[u(x)]$ for different values of c , ordinary product.

interval $(0, 1)$. In Figure 4, we plot the solution variance at point $x = 0.5$ as a function of c . It appears that variances corresponding to ordinary and Wick products are close for small c , but diverge quickly as c increases: in the same range of the values of c , the variance of the solution increases from $O(10^{-5})$ to $O(10^6)$ for the ordinary product, and from $O(10^{-5})$ to $O(10^1)$ for the Wick product.

In summary, although the solutions of equation (5.1) with Wick and ordinary products belong to $L_2(\mathbb{F}; H_0^1(0, 1))$, the two different products lead to many differences in the properties of the solution. The differences are especially pronounced when the noise K is large. A more detailed analysis shows that many numerical characteristics of the solutions differ by

FIGURE 3. Variance of the solution for different values of c .

Left: Ordinary product; Right: Wick product

FIGURE 4. Variance of the solution at $x = 0.5$, as a function of c .

the factor of e^{c^2} , which is related to the equality $e^{c\xi - \frac{1}{2}c^2} \diamond e^{-c\xi - \frac{1}{2}c^2} = 1$. At this point, we make the statement intentionally vague, because this scaling by the factor e^{c^2} seems to have many important ramifications and must be studied further, both analytically and numerically.

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